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Vibrational Excess Entropy of Equimolar KCl-KBr Solid Solution

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The Debye-Waller factors of the ions in equimolar KCl-KBr solid solution at 300 K have been determined from X-ray diffraction data. It was observed that the amplitudes of thermal vibration are slightly higher in the solid solution than in the pure components. The corresponding vibrational excess entropy was found to be about 10% of the entropy of ideal

In solid solutions, the differences in the sizes of the component atoms give rise to small random displacements from the sites of the average lattice. One effect which arises from these static displacements is that the intensities of the X-ray reflexions are reduced by an exponential factor, which is similar to the Debye-Waller (D-W) factor produced by the thermal displacements. In addition to the existence of the size effects, the vibrational states of the constituent atoms may change in the solid solution formation as a result of new interatomic interactions. If these changes occur, the entropy of the solid solution will deviate from that of ideal mixing. Since it is well known for metallic alloys that the vibrational contribution is crucial in accounting for deviations from the entropy of ideal mixing 1, 2, the question arises of whether this contribution is also of importance in alkali halide solid solutions. In the present work, the vibrational excess entropy of equimolar KCl-KBr solid solution was determined from X-ray diffraction data. The data of diffuse X-ray scattering indicate 3 that no departures from random distribution appear in this system.

The component salts utilized in preparation of the solid solution were of quality "supra pure", supplied by E. Merck A. G. (Darmstadt, Germany). The mixture, consisting of 50 mole per cent of both components, was melted in an evacuated quartz ampulla at 800 °C. The solid solution was annealed at 600 °C for 2 h for homogenization, and subsequently ground in a ball-mill at room temperature until the average particle size was about 4 μ m.

The integrated intensities of the Bragg reflexions from three samples, which were prepared in a mould under uniaxial pressure, were measured on the same relative scale using monochromatized CuKa radiation. The experimental procedure was essentially similar to that described by INKINEN et al. 4. After the measured intensities had been corrected for preferred orientation 5 and for thermal diffuse scattering 6, they were

¹ O. J. Kleppa, in: Metallic Solid Solutions (Edited by J. FRIEDEL and A. GUINIER), Benjamin, New York 1963. M. SIMERSKA, Acta Met. 13, 1153 [1963].

converted to the experimental structure factors, $F_{\rm exp}$. The theoretical structure factors, $F_{
m theor}$, were calculated from relativistic Hartree-Fock free-ion form factors 7. The scale factor of the values $F_{\rm exp}$ and the D-W parameters, $B_{\mathbf{K}}^{\mathbf{ss}}$ and $B_{\mathbf{Cl},\mathbf{Br}}^{\mathbf{ss}}$, of the ions in solid solution were determined on the assumption that the values $F_{\rm exp}$ and $F_{\rm theor}$ coincided at high values of $\sin \theta/\lambda$, by employing the difference-Fourier method. The dispersion-corrected 8 values of $F_{\rm exp}$ are listed in Table 1, together with the values $F_{\rm theor}$ and the differences $\Delta F = F_{\rm exp} - F_{\rm theor}$. The limits of error of the values $F_{\rm exp}$ do not include the estimated uncertainty of 1.5% in the scale factor.

hkl	$F_{ m exp}$	$F_{ m theor}$	ΔF
111	27.93 ± 0.30	27.07	0.86
200	135.95 ± 1.25	139.11	-3.16
220	114.70 ± 0.60	114.93	-0.23
311	20.84 ± 0.25	21.14	-0.30
222	96.32 ± 0.65	98.30	-1.98
400	87.29 ± 0.65	85.93	1.36
331	19.69 ± 0.45	18.97	0.72
420	77.47 ± 0.75	76.24	1.23
422	67.82 ± 0.70	68.36	-0.54
333, 511	16.22 ± 0.45	16.87	-0.65
440	55.54 ± 0.75	56.11	-0.57
531	14.97 ± 0.30	14.73	0.24
442,600	51.90 ± 0.60	51.18	-0.72
620	47.34 ± 0.50	46.81	0.53
622	42.34 ± 0.50	42.90	-0.56
444	37.94 ± 0.75	39.38	-1.44
551, 711	10.85 ± 0.30	10.95	-0.10
640	36.40 ± 0.50	36.20	0.20
642	33.17 ± 0.45	33.32	-0.15
800	28.83 ± 0.55	28.31	0.52

Table 1. Experimental dispersion-corrected structure factors F_{exp} of equimolar KCl-KBr solid solution at 300 K, theoretical structure factors $F_{
m theor}$ calculated from relativistic Hartree-Fock free-ion scattering factors with Debye-Waller parameters evaluated in the present work, and differences $\Delta F = F_{\rm exp} - F_{\rm theor}$.

The D-W parameters arrived at in the present work are indicated in Table 2, together with the previous experimental and theoretical results for the ions in KCl and KBr. Since the spread in the B values of the component salts is comparable to the average uncertainty interval (±0.10 Å2) of the experimental data, justification exists for regarding their mean values as the D-W parameters of the pure components.

Comparison of the D-W parameters of the ions in solid solution with the values obtained by averaging the results for KCl and KBr discloses that the values of the positive and negative ions are increased in the



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P. Luova, private communication.

O. INKINEN, A. PESONEN, and T. PAAKKARI, Ann. Acad. Sci. Fenn. A VI No. 344 [1970].

⁵ M. JÄRVINEN, M. MERISALO, A. PESONEN, and O. INKINEN, J. Appl. Cryst. (in press).

⁶ D. R. CHIPMAN and A. PASKIN, J. Appl. Phys. 30, 1998

⁷ P. A. Doyle and P. S. Turner, Acta Cryst. A 24, 390

⁸ D. T. CROMER, Acta Cryst. 18, 17 [1965].

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Crystal	Debye-Waller parameter		
KCl-KBr	$B_{\rm K}^{\rm ss} = 2.71 \pm 0.10 \text{ a}$	$B_{\text{Cl,Br}}^{\text{ss}} = 2.26 \pm 0.10 \text{ a}$	
KCl	$B_{\rm K} = 2.08 \pm 0.05 \text{ b}$ 1.93 c	$B_{\rm Cl} = 2.06 \pm 0.05 \text{ b}$ 1.99 c	
KBr	$B_{ m K} = 2.45 \pm 0.15 { m d} \ 2.20 \pm 0.10 { m e} \ 2.21 { m c}$	$B_{\mathrm{Br}} = 2.23 \pm 0.15 \mathrm{d} \ 2.28 \pm 0.10 \mathrm{e} \ 2.24 \mathrm{c}$	

a Present work.

b Ref. 9: Observed; polycrystal, X-rays.

c Ref. 10: Calculated from the normal mode data.

d Ref. 11: Observed; polycrystal, X-rays.

e Ref. 12: Observed; single crystal, neutrons.

Table 2. Debye-Waller parameters B (in Å²) of the ions in the equimolar KCl-KBr solid solution and in its pure components at 300 K.

solid solution formation by $\Delta B_{\rm K} = 0.54~{\rm Å}^2$ and $\Delta B_{\rm Cl, Br} = 0.10~{\rm Å}^2$, respectively. The large change in the value of the K ion, as compared with the small increase in $B_{\rm Cl, Br}$, gives an indication of a local static distortion arising from different sizes of the Cl and Br ions. The size-effect displacement of the K ion is 13 $s=0.0766~{\rm Å}$, so that $\Delta B_{\rm K}({\rm static})=8~\pi^2~s^2=0.46~{\rm Å}^2$. Subtraction of this value from the results obtained for $B_{\rm K}^{\rm ss}$ and $\Delta B_{\rm K}$ gives

 B_{κ}^{ss} (thermal) = 2.25 Å² and $\Delta B_{\rm K}$ (thermal) = 0.08 Å². The latter value is comparable to the difference $\Delta B_{\rm Cl,\,Br}$ mentioned above. The results obtained have thus indicated that the amplitudes of thermal vibration are slightly higher in the actual solid solution than in the ideal one, which implies slightly looser binding and higher vibrational entropy in the former.

The entropy ΔS of the solid solution is composed of the entropy of ideal mixing and of additional contributions, the number of which will depend upon the physical properties of the system under consideration. For metals, additional entropy terms attributable to differences in electronegativity and valence contribute ¹ to ΔS , but may be ignored for alkali halide solid solu-

tions. Since the electrical conductivity measurements of Ambrose and Wallace ¹⁴ gave no indication of an abnormal population of vacancies in KCl-KBr solid solution, the entropy contribution attributable to the existence of vacancies is insignificant in this system.

The above considerations lead to the assumption that ΔS is composed of the entropy of ideal mixing and of the vibrational contribution: $\Delta S = \Delta S^{\mathrm{id}} + \Delta S^{\mathrm{vib}}$. For the equimolar solid solution, $\Delta S^{\mathrm{id}} = 5.76$ J/deg. mole. The molar vibrational entropy was calculated from equation ²

$$\Delta S^{\text{vib}} = \frac{3}{2} R \ln (\overline{B^{\text{ss}}}/\overline{B})$$
, (1)

where R is the gas constant, \overline{B}^{ss} the average \overline{D} -W temperature parameter of the solid solution, and \overline{B} the corresponding quantity calculated from the B values of the pure components. By application of the averaging scheme 2 , which weights the B values of the individual components in proportion to both their concentration and the number of electrons associated with each component, there is obtained from Eq. (1)

$$\Delta S^{\mathrm{vib}} = 0.5 \pm 0.2 \mathrm{\ J/deg.\ mole}$$
 .

This result is somewhat less than the value $\Delta S^{\mathrm{vib}} = 0.84 \,\mathrm{J/deg.}$ mole calculated 15 on the basis of the Debye temperatures derived from experimental elastic-constant data. With the value of ΔS^{vib} obtained in this work, the total entropy will be $\Delta S = 6.3 \,\mathrm{J/deg.}$ mole. This result coincides with that deduced 16 from data of the experimental heat of formation and activity. It thus appears safe to conclude that the formation of an equimolar KCl-KBr solid solution is characterized by a positive vibrational excess entropy, which is about 10% of ΔS^{id} . This result is at variance with one of the basic ideas of Wasastierna's theory 17 , that the vibrational contribution to ΔS is zero.

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⁹ L. K. PATOMÄKI and M. V. LINKOAHO, Acta Cryst. A 25, 304 [1969].

¹⁰ W. J. L. Buyers and T. Smith, J. Phys. Chem. Solids 29, 1051 [1968].

¹¹ V. Meisalo and O. Inkinen, Acta Cryst. 22, 58 [1967].

¹² A. W. PRYOR, Acta Cryst. 20, 138 [1966].

¹³ J. HIETALA, Ann. Acad. Sci. Fenn. A VI No. 122 [1963].

¹⁴ J. E. Ambrose and W. E. Wallace, J. Phys. Chem. **63**, 1536 [1959].

¹⁵ D. L. FANCHER and G. R. BARSCH, J. Phys. Chem. Solids 30, 2517 [1969].

¹⁶ W. H. McCoy and W. E. WALLACE, J. Amer. Chem. Soc. 78, 5995 [1956].

¹⁷ J. A. WASASTJERNA, Soc. Sci. Fenn., Comm. Phys.-Math. XV No. 3 [1949].